

## Glaciers of the McMurdo dry valleys: Terrestrial analog for Martian polar sublimation

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[1] The surfaces of the Martian north and south polar residual caps are marked by unusual ice features: Dark spiralesque troughs up to 1 km deep, 10 km wide, and 300 km long appear on both ice caps, and circular pits that make up the “Swiss cheese” terrain appear on the south polar cap. Both types of features are of interest to researchers as a potential means of understanding ice composition and flow rates. Some glaciers of the McMurdo dry valleys have surface features unknown elsewhere on terrestrial glaciers, including canyons over 6 km long, 100 m wide, and tens of meters deep and basins up to 100 m across. High sublimation, dust accumulation, and very little melting is key to their origin. These processes and ice landforms are suggested as terrestrial analogs for the sublimation behavior of Martian ice caps, where dust accumulation and sublimation are significant but surface melting is absent. We have developed a solar radiation model of canyon formation and have applied it to the Martian polar caps. The modeled processes do well to describe direct and reflected radiation within V grooves, a process that may be significant in the development of the spiral troughs and Swiss cheese terrain. The model fails to reproduce the low observed slopes of the Martian troughs. The grooves are too shallow, with opening angles of  $\sim 165^\circ$  compared with model predictions of  $\sim 90^\circ$ . The reason for the failure may be that we have not included creep closure, which should flatten their slopes.

*INDEX TERMS:* 5462 Planetology: Solid Surface Planets: Polar regions; 1827 Hydrology: Glaciology (1863); *KEYWORDS:* dry valleys, Mars, glaciers, sublimation, Swiss cheese terrain

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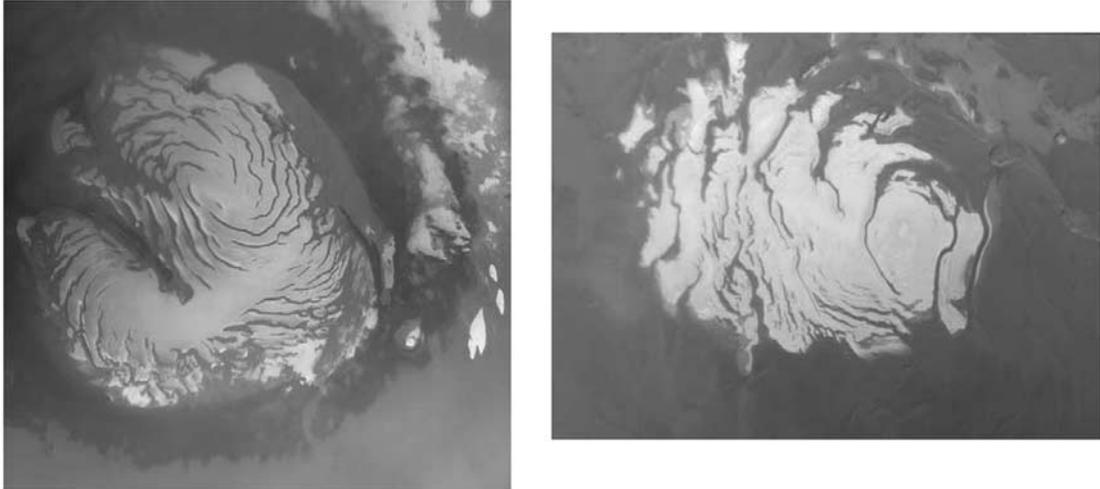
### 1. Introduction

[2] The obvious summer features on the residual polar caps of Mars are dark spiralesque markings (Figure 1). These markings are dust-mantled troughs developed in the ice. Troughs are up to 1 km deep and 10 km wide; in the south polar cap some are 150 km long, and in the north polar cap they can be up to 300 km long. Seen at high resolution, the trough walls (especially the equator-facing wall) are step-like outcrops of quasi-horizontal strata or layers of alternating clean ice and dusty ice or ice-bonded dust layers. Exposure of stratification clearly indicates that these canyons have been eroded. Evidence points to deeper erosion of the canyons near the edges of the polar caps [Ivanov and Muhleman, 2000b].

[3] There is much discussion, but no consensus, concerning the origin of Martian polar troughs. Several current theories are being explored. The first is that the troughs originate near the edges of the polar deposits and migrate toward the pole by preferential sublimation of ice from the steep, equator-facing side and accumulation on the pole-facing side, as shown in Figure 2 [Howard, 1978; Howard et al., 1982; Fisher, 1993]. Ivanov and Muhleman [2000a]

used a sublimation model with no flow to reproduce, relatively successfully, the general shape of the troughs over a time span of 3–16 million years. The key was to provide an initial albedo contrast on the ice cap surface. Fisher [1993, 2000] combines the idea of trough migration with ice cap flow outward from the center; troughs migrate toward the center of the ice dome while ice flows outward from the center until the two rates are balanced. This model suggests that the spiral pattern of troughs and valleys results from the nonsymmetric distribution of the ice flow centers and associated velocity fields around the pole. In addition, there have been a number of recent attempts to reproduce the troughs through sublimation and accumulation modeling. Some models which incorporate flow [e.g., Hvidberg, 2000; Larsen, 2000] tend to have trouble keeping the troughs open over time, possibly owing to the ice rheology used for modeling. There has also been some exploration of the role of eolian erosion in maintaining the troughs [Howard, 2000].

[4] The structural geology, geomorphology, and ice processes of individual, specific troughs are relatively unexplored. There are few published mechanisms for their initiation, and estimates of growth rates are sparse and generally only from rough parameterizations of accumulation and sublimation rates [Ivanov and Muhleman, 2000b]. Existing theories for trough behavior are limited. Current



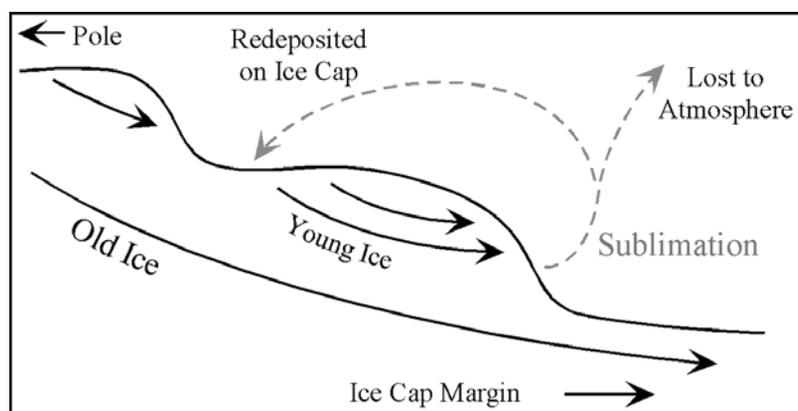
**Figure 1.** Mars Orbiter Camera (MOC) image of the Martian (left) north (MGS MOC Release MOC2-231, 22 May 2000) and (right) south (MGS MOC Release MOC2-225, 27 April 2000) polar regions during their respective summers. The central white area in each image is the residual cap. The north polar residual cap is  $\sim 1100$  km across and is cut by dark, spiral troughs; the south polar residual cap is  $\sim 420$  km across. MOC images were obtained through the Malin Space Science Systems website ([http://www.msss.com/moc\\_gallery](http://www.msss.com/moc_gallery)).

theories require the troughs to be relatively recent developments, or else some as yet unexplored mechanism is required for control of their growth or for trough removal from the ice caps. None of the proposed models work as steady state solutions where the current trough configuration is one stage in an ongoing process. Instead, the troughs need to be relative newcomers on the scene, either developed at the ice cap edge and in the process of growing inward or recently etched into the surface via sublimation.

[5] The goal of this paper is to use the basins and canyons found on the Canada and Taylor Glaciers in the McMurdo dry valleys, Antarctica, as terrestrial analogs for understanding the spiral troughs of the Martian north polar cap and the Swiss cheese terrain of the Martian south polar cap. In the dry valleys, basins and canyons are formed within the ablation zones of many of the glaciers as initially small features that grow as they are carried to the glacier margin by ice flow. We propose a similar evolution on Mars starting high up on the ice cap. We hypothesize that the growth of

the spiral troughs seen on the Martian north polar cap is initiated by locally enhanced sublimation via a perturbation in surface slope and aspect or through heterogeneous albedo due to windblown dust or grain annealing. Evidence for both small trough features and for dust deposits near the centers of the ice caps can be seen in the Mars Orbiter Camera (MOC) imagery and Mars Orbiter Laser Altimeter (MOLA) data [Howard, 2000]. Once established, the features would grow via enhanced sublimation rates on their steep, equator-facing sides, a mechanism similar to that proposed in the existing Martian scarps models and similar to the behavior of dry valley basins and canyons. In support of this hypothesis we present a comparison of the dry valley and Martian ice cap environments, a discussion of the morphology and development of canyons in the dry valleys, and a discussion of the Martian features.

[6] Our hypothesis assumes that sublimation associated with enhanced solar absorption is the principal means by which grooves are excavated and maintained against closure



**Figure 2.** Cartoon representation of sublimation and accumulation on a scarp on the Martian ice cap [adapted from Fisher, 1993].

by ice inflow. To examine this assumption, a simple V groove model, developed by *Pfeffer and Bretherton* [1987] to model radiation absorption within crevasses, is applied to grooves on the Martian north polar residual cap. The model calculates both incident and reflected shortwave radiation within a simplified V groove for a specified solar azimuth. Results from this work will contribute to determining ice composition, since groove stability requires that sublimation roughly balances ice inflow. Knowledge of the composition of the Martian polar caps will help determine the history of volatile substances (water and carbon dioxide) in the Martian atmosphere and on the Martian surface.

[7] In addition to the trough features seen on both Martian ice caps, MOC has recently returned images of strange features on the south polar residual cap that have been termed the “Swiss cheese” terrain [*Sullivan et al.*, 2001; *Malin et al.*, 2001] (Figure 3). *Byrne and Ingersoll* [2000] have attributed the origin of this terrain to ablation via sublimation of CO<sub>2</sub> ice. The circularity of pits is probably a result of sublimation at high southern latitude where solar elevation is nearly constant throughout the day and where therefore ablation on the walls of depressions is roughly radially symmetric. The formation of the Swiss cheese pits takes place on century timescales, the same timescales as the formation of canyons and basins in Antarctic ice glaciers. Significant interannual changes have been observed [*Malin et al.*, 2001]. The morphologies present within the Swiss cheese terrains are quite varied and point out the possibility that they consist of layers of multiple ices (H<sub>2</sub>O, CO<sub>2</sub>, and CO<sub>2</sub> clathrate are all likely). Less volatile ice layers, H<sub>2</sub>O ice perhaps, may form mesa tops and hole bottoms, with more volatile ice layers in between, which can rapidly ablate outward from an initial point, as is postulated by *Thomas et al.* [2000].

[8] As with the polar troughs, further insight into the working of these features can be gleaned from comparison with the basins seen on dry valley glaciers. We propose that these features are the result of sublimation of the surface in the absence of effective wind removal of debris. Sheltered conditions within the basins may aid retention of debris, as seen in the dry valley basins. The first areas of the Swiss cheese terrain to become free of winter dry ice snows, leaving dusty, annealed layers of dry ice behind, are scarps, as expected. The snow-free scarps should have a large effect on the local radiative environment, as may be evidenced by the moats commonly seen fringing the foot of most scarps, though other possible explanations for the moats have also been put forward [*Thomas et al.*, 2000].

## 2. Physical Environment: Antarctic Dry Valleys Versus Martian Polar Regions

[9] Though very different environments, there are enough critical similarities between the Antarctic dry valleys and the Martian polar regions to allow for fruitful comparison between the two. Both are cold, dry regions where energy for ablation is limited, accumulation and sublimation rates are relatively small, and ice flow speeds are low. In these regions, strange topography can be etched into ice surfaces as a result of dominance of sublimation and differential ablation rates.

[10] Table 1 presents a number of physical environmental parameters for both regions. With the exception of accu-



**Figure 3.** Martian south polar residual cap “Swiss cheese” terrain (MGS MOC Release MOC2-211, 8 March 2000). The image is illuminated by sunlight from the upper left. This image shows an upper layer of the Martian south polar residual cap that has been eroded, leaving flat-topped mesas into which are set circular depressions. The tallest mesas shown here stand  $\sim 4$  m high. Terrain such as this is found only on the south polar residual cap, leading to speculation that these landforms may result from the CO<sub>2</sub> ice thought to be common in the south polar region. The image covers an area  $3 \times 9$  km near  $85.6^\circ\text{S}$ ,  $74.4^\circ\text{W}$  at a resolution of  $7.3 \text{ m pixel}^{-1}$ . This picture was taken during early southern spring on 3 August 1999.

**Table 1.** Comparison of Physical Environmental Parameters for the Antarctic Dry Valleys and Martian Polar Caps

	Antarctic Dry Valleys	Martian North Polar Cap	Sources for Martian Data
Gravity, $\text{ms}^{-2}$	9.81	3.72	
Solar constant (top of atmosphere), $\text{W m}^{-2}$	1380	594	
Atmospheric pressure, mbar	1000	5–20	
Atmospheric water vapor pressure, mbar	0–9	$\sim 10^{-3}$	
Accumulation rates of $\text{CO}_2$ ice, $\text{m yr}^{-1}$		1–2	<i>Smith et al.</i> [2001]
Accumulation rates of $\text{H}_2\text{O}$ ice, m water equivalent $\text{yr}^{-1}$	$10^{-2}$ – $10^{-1}$	$10^{-4}$	<i>Clifford et al.</i> [2000]
Sublimation rates of $\text{CO}_2$ ice, $\text{m yr}^{-1}$		$1$ – $2^a$	
Sublimation rates of $\text{H}_2\text{O}$ ice, m water equivalent $\text{yr}^{-1}$	$10^{-1}$ <sup>b</sup>	$10^{-4}$ – $10^{-3}$	<i>Jakosky and Haberle</i> [1992]
Melt rates, m water equivalent $\text{yr}^{-1}$	$10^{-1}$ – $10^0$ <sup>b</sup>	none	
Mean annual surface temperature, K	$\sim 255$	$\sim 155$	<i>Clifford et al.</i> [2000]
Flow speeds, $\text{m yr}^{-1}$	$10^0$ – $10^1$	order $10^{-2}$	<i>Fisher</i> [2000]
Dust deposition rates, $\text{gm}^{-2} \text{yr}^{-1}$	0.5–1	20–200	<i>Thomas et al.</i> [1992]
Volumetric fraction of debris in ice	$\leq 1\%$	1–50+%	<i>Clifford et al.</i> [2000]

<sup>a</sup>All over-winter accumulation sublimates.

<sup>b</sup>In ablation zone.

mulation and ablation of  $\text{CO}_2$  ice, which does not exist under ambient surface conditions on Earth, polar processes on Mars occur more slowly than on Earth. Atmospheric pressure and water vapor pressure are several orders of magnitude smaller; accumulation and ablation rates for  $\text{H}_2\text{O}$  ice are several orders of magnitude smaller; the ice is much colder, resulting in a much smaller range of probable flow speeds. Gravity is about a third that of Earth, and solar radiation is a little less than half that of Earth. Length of day and obliquity are among the few physical parameters that are nearly equal. However, the ratios of accumulation versus sublimation for  $\text{H}_2\text{O}$  ice are similar: Sublimation is up to 1 order of magnitude greater than accumulation. Similarly, the ratio of accumulation to flow speed is  $\sim 10^2$ , and the ratio of sublimation to flow speed is  $\sim 10^1$ – $10^2$  (see Table 1). This similarity may be part of the key to understanding why large roughness features form both in the McMurdo dry valleys and on Mars in spite of environmental differences.

[11] Controls on the rate of sublimation are somewhat different on Earth and Mars. In the Earth atmosphere, sublimation is predominantly a function of the vapor pressure difference between the surface and the air and of the speed of turbulent dispersion of the air. Consequently, equations for sublimation incorporate the vapor pressure gradient, wind speed (as a measure of turbulence), and buoyancy effects. The latent energy flux  $Q_E$  for Earth conditions can be written as

$$Q_E = -u^* k_o z \left( \frac{0.622 \rho_i}{P} \right) \frac{\partial e}{\partial z} L_s, \quad (1)$$

where  $u^*$  is a characteristic scale of wind speed,  $k_o$  is the von Karman constant,  $z$  is the height at which wind speed and vapor pressure are measured,  $\rho_i$  is the ice density,  $P$  is atmospheric pressure,  $\partial e/\partial z$  is the gradient of atmospheric vapor pressure with height, and  $L_s$  is the latent heat of sublimation [Paterson, 1994]. The variable  $u^*$  is written as

$$u^* = \frac{uk_o}{\ln \left( \frac{z}{z_o} \right)}, \quad (2)$$

where  $u$  is the average wind speed and  $z_o$  is the surface roughness parameter. Since sublimation on Earth is driven by vapor pressure gradients and wind, only in extreme conditions will sublimation be limited by energy availability. (Sublimation cools the surface until the change in

vapor pressure gradient reduces sublimation rates to match energy availability.) This type of energy limitation of sublimation does not occur in the dry valleys during the summer. Consequently, for our purposes here, sublimation on Earth can be considered to be entirely forced by wind and vapor pressure gradients.

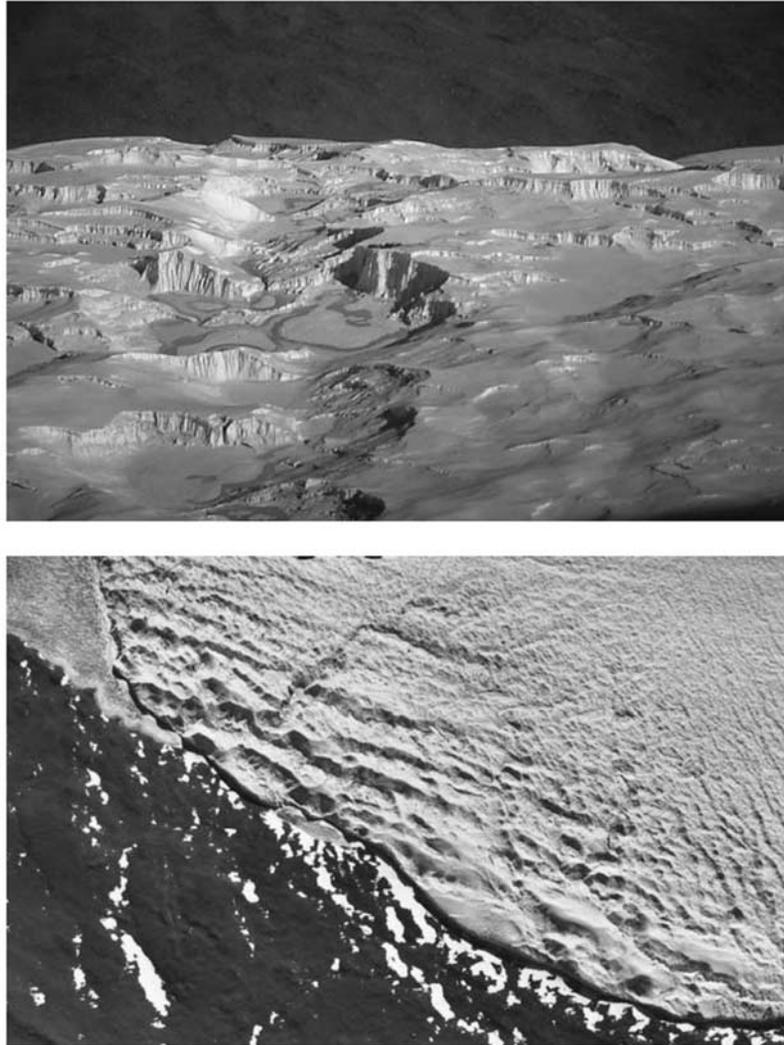
[12] For Mars, owing to the low atmospheric pressure ( $< 20$  mbar), the dominant process controlling water vapor flux is natural convection. Water vapor is more buoyant in the Martian  $\text{CO}_2$  atmosphere than it is on Earth [Toon *et al.*, 1980]. At the same time, forced convection via wind and vapor pressure gradients plays a relatively small role. For Martian conditions the vapor flux is written with two equations [Toon *et al.*, 1980], the first for turbulent transfer as a result of winds and the second for turbulent transfer as a result of natural convection:

$$Q_{E1} = 0.002 \rho_w u L_s \quad (3)$$

$$Q_{E2} = 0.017 \rho_w D \left[ \left( \frac{\Delta \rho}{P} \right) \left( \frac{g}{\nu^2} \right) \right]^{\frac{1}{2}} L_s \quad (4)$$

In these equations,  $\rho_w$  is the water vapor density,  $D$  is a diffusion coefficient for  $\text{H}_2\text{O}$  in a  $\text{CO}_2$  atmosphere,  $P$  is atmospheric pressure,  $\Delta \rho/P$  relates water vapor pressure to atmospheric pressure,  $g$  is gravity, and  $\nu$  is the viscosity of  $\text{CO}_2$ . Values for the constants are given by Toon *et al.* [1980]. These equations are very similar to the sublimation equations used by others [e.g., Ingersoll, 1970; Haberle and Jakosky, 1990]. Of these two terms,  $Q_{E2}$  generally dominates over  $Q_{E1}$ . This means the primary control on sublimation lies in  $Q_{E2}$  in the  $\Delta \rho/\rho$  term, which is governed by surface temperature. Surface temperature is governed primarily by energy absorption. Thus Martian sublimation is driven primarily by energy absorption at the surface, which implies that slope, aspect, and albedo enhancement of energy absorption will directly increase ablation.

[13] An important similarity between the dry valley glaciers and the Martian polar caps is that both regions are energy-limited. In the dry valleys, temperatures are typically below freezing, and much of the ablation is due to sublimation resulting from the persistent high winds, which greatly reduce or eliminate energy available for melt



**Figure 4.** Aerial photographs of the lower Canada Glacier. (top) Oblique view of several of the roughly circular basins. This photo was taken looking southwest; the asymmetry of the equator-facing vertical wall and pole-facing, shallow slope can be seen. (bottom) Larger section of the glacier toe. The circular basins can be seen best in the lower right section of the image.

[Lewis *et al.*, 1998]. Therefore if energy absorption is enhanced on a surface, for example, via slope, aspect, or albedo, that energy will be primarily converted to melt, and the surface will experience significantly more ablation. For the same amount of energy, melt is  $\sim 10$  times more efficient than sublimation at removing ice, so only a small amount of “extra” energy devoted to melt is needed to make a large impact on local ablation.

[14] On the Martian polar caps, energy availability limits sublimation much as energy availability limits melt in the dry valleys. Temperatures do not reach the melting point on the Martian polar caps. The altitude of the south polar cap is such that barometric pressure is less than the triple-point pressure, so even with a remarkable and unlikely coincidence of conditions that could allow the melting point to be attained, the ice would first sublimate. Local slope, aspect, and albedo do have dramatic effects, enhancing local sublimation rates on the Martian polar caps much the same as they allow increased melt in the dry valleys. When coupled with the longer timescales over which features form on Mars,

enhancements in sublimation are enough to form sublimation pits or canyons. On Earth it is the factor of 10 difference between melt and sublimation that allows canyon and pit formation. Through increased melt the dry valley channels on Taylor Glacier, which are roughly  $10 \text{ m} \times 100 \text{ m} \times 1 \text{ km}$  in size, form in  $\sim 1000$  years (see section 3). On Mars, features  $1 \text{ km} \times 10 \text{ km} \times 100 \text{ km}$  may form over an estimated time span of  $5 \times 10^7$  years (this assumes an ice flow of  $10^{-2} \text{ m yr}^{-1}$  for a distance of 500 km, the radius of the ice cap). So we have a factor of  $5 \times 10^4$  more time on Mars to form features 100 times larger. This implies that processes can be up to 500 times slower on Mars than on Earth and still produce troughs of the sizes seen in the residual polar caps.

### 3. Glacier Surface Morphology in the McMurdo Dry Valleys, Antarctica

[15] The basins and canyons found on the dry valley glaciers appear to form as a result of three principal environmental factors.



**Figure 5.** Lower Taylor Glacier, Taylor Valley, Antarctica. This image shows how the canyons at the glacier toe appear to be a continuation of banding in the ice seen far up glacier. The section of glacier shown here is  $\sim 6$  km long. North is toward the top of the image.

[16] 1. The glaciers are frozen to their beds and move comparatively slowly. Features in the ablation zone of the glaciers have hundreds to thousands of years to develop.

[17] 2. The environment is dry and windy, and deposits of dust and sand (either through eolian processes, medial moraines, or rock avalanches) are ubiquitous on the glaciers ( $0.5\text{--}1\text{ g m}^{-2}\text{ yr}^{-1}$ ). The debris melts into the ice, if only a short distance. Generally, meltwater is absent or insufficient to wash the sand and dust off the glacier.

[18] 3. The environment is energy poor. Temperatures are typically below freezing, and there is relatively little incoming energy,  $\sim 40\text{ W m}^{-2}$  on average over the summer season [Lewis *et al.*, 1998]. The low levels of available energy, coupled with the cold, windy conditions, results in very little melt on the glacier surfaces. What little energy is available goes primarily to sublimation. As a result, small changes in energy receipt on the glacier surface as a result of albedo, slope, and aspect greatly affect melt production [Lewis *et al.*, 1998].

[19] The basins on the dry valley glaciers appear to nucleate around debris. Isolated patches of debris form  $\sim 30$ -cm-deep cylindrical holes, “cryoconite holes,” in the glacier surface [Mueller *et al.*, 2001; Wharton *et al.*, 1985]. These holes vary in radius from a few centimeters to over a meter. At small radii the holes retain an ice cover and appear to grow very slowly. At radii  $>0.5$  m the holes begin to lose their ice covers. At this point, if debris and meltwater within the holes can drain, the holes can rapidly deepen and widen (up to  $\sim 0.50$  m per summer in both depth and width [Lewis, 2002]). In the process the holes take on an asymmetric geometry. The pole-facing wall develops into a shallow pole-facing slope because of shading, while the equator-facing wall remains vertical and melts back relatively rapidly (Figure 4). A similar process applies to the long channels or canyons, such as on Taylor Glacier, but the source of debris is medial moraines that originate at the confluences of the glacier tributaries (Figure 5). Basin and valley formation therefore requires both the initial presence of debris and the removal of debris and meltwater when the feature reaches a critical size and results in a feature

geometry that enhances solar radiation absorption in some locations and reduces solar absorption in others.

[20] The Taylor Glacier canyons (Figures 5 and 6 (right)) appear to start  $\sim 3$  km from the glacier terminus as  $\sim 1$ -m-deep grooves. These grooves, initiated by medial moraine debris, grow as they are carried down the glacier by ice flow. Ice speed is  $4\text{--}7\text{ m yr}^{-1}$ , so the 3-km distance from the point of canyon initiation to the glacier terminus represents 400–750 years of evolution. At the terminus the grooves have become 10- to 20-m-deep chasms in the glacier surface, formed by differential ablation rates and modified by water flowing within the canyon. The canyon width initially varies, with wide pool-like features separated by relatively narrow passages. At some point these features vanish, and the width to depth reaches a nearly constant value.

[21] The Canada Glacier basins are more randomly distributed than the canyons on Taylor Glacier because they nucleate around debris patches [Lewis, 2002]. Owing to their more random distribution the basins seem to function more as isolated features, each growing at its own rate. Timescales for basin growth on Canada Glacier are 400–3000 years (ice flow rates of  $1\text{--}8\text{ m yr}^{-1}$  over a distance of up to 3 km).

#### 4. Martian Ice Caps: Their Climate and Behavior

[22] Current consensus is that the composition of the Mars north polar residual cap is  $\text{H}_2\text{O}$  ice, with only a seasonal  $\text{CO}_2$  frost that completely sublimates during the summer [Clifford *et al.*, 2000]. However, there are alternate possibilities that would include  $\text{CO}_2$  clathrate and/or dry ice at depth [Mellon, 1996; Ross and Kargel, 1998; Kargel and Lunine, 1998; Kargel *et al.*, 2000]. The composition of the south polar cap is less well understood, but accumulating evidence points toward a form of  $\text{CO}_2/\text{H}_2\text{O}$  clathrate or some layering of  $\text{CO}_2$ ,  $\text{H}_2\text{O}$ , and  $\text{CO}_2$  clathrate ices [Jakosky *et al.*, 1995; Kargel and Lunine, 1998; Thomas *et al.*, 2000]. On both caps, air and surface temperatures, at the current obliquity, do not reach or come close to attaining the melting point of  $\text{H}_2\text{O}$ . Winter temperatures on both caps



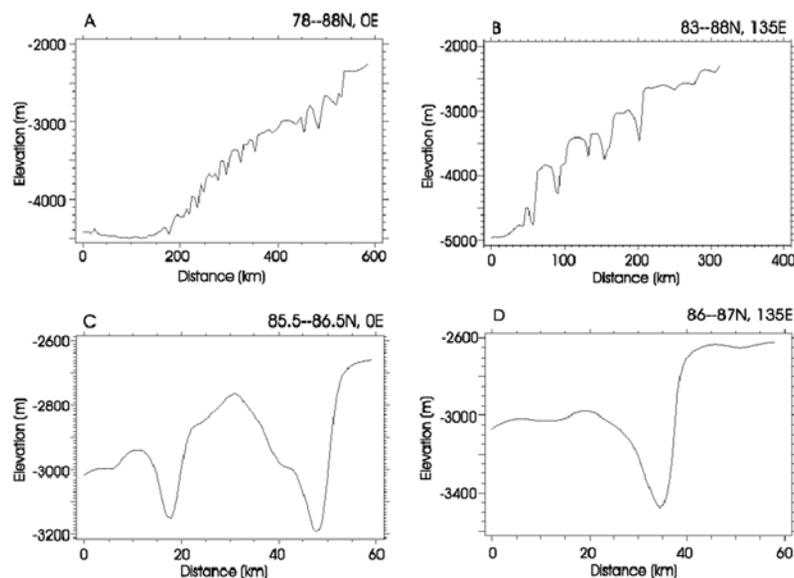
**Figure 6.** (left) Canada and (right) Taylor Glaciers, Taylor Valley, Antarctica. The large roughness features can be seen on the lower reaches of the glaciers. For scale, both glaciers are  $\sim 3$  km wide.

drop to 148 K, allowing the deposition of  $\text{CO}_2$  frost. Summer high temperatures for the north cap and outlying south polar layered deposits are  $\sim 205$  K [Kieffer *et al.*, 1976; Kieffer and Zent, 1992; Jakosky *et al.*, 2000], resulting in the sublimation of the winter accumulation of  $\text{CO}_2$  frost as well as some portion of the  $\text{H}_2\text{O}$  ice of the cap. This has been seen in the Mars Atmospheric Water Detection data [Farmer and Doms, 1979], which shows enhanced water vapor over the north pole in summer. On the south polar cap, summer temperatures do not always rise measurably above the dry ice frost point ( $\sim 148$  K), so that  $\text{CO}_2$  frost, deposited during the fall and winter, may not completely sublimate during the summer season [Kieffer, 1979; Jakosky and Haberle, 1992]. The distinct morphologies of the two polar caps also support a different surface ice composition. Since temperatures stay well below the melting point for both caps, ablation will only occur via sublimation and wind erosion. Since sublimation increases with temperature [Jakosky *et al.*, 1993; Ivanov and Muhle-

man, 2000a], ice slopes that face the Sun or are dust-covered will warm and sublimate faster. This results in the formation of accentuated slopes or depressions that receive yet more energy, leading to a positive feedback.

[23] Sublimation is not the only process controlling trough geometry. Because the ice cap is so thick and because the troughs are so deep, trough location and shape will be affected by ice thickness and rheology, and trough closure by ice flow will probably be significant. The north polar cap rises  $\sim 3$  km over the surrounding terrain. Zuber *et al.* [1988] calculated an ice cap thickness of  $\sim 3.5$ –4 km. For the south ice cap, P. Schenk and J. Moore (see <http://www.lpi.usra.edu/research/msp/msp.html>) estimate the south polar cap to be a broad convex dome  $\sim 500$  km in diameter with a maximum height of 3 km above the surrounding plains.

[24] Such thick ice will flow [Budd *et al.*, 1986; Greve, 2000]. If the caps were at the measured surface temperatures throughout, they would flow very slowly, but the ice is



**Figure 7.** Cross-section profiles of troughs on the Martian north polar cap, derived from Mars Orbiter Laser Altimeter (MOLA) data: (a)  $78^{\circ}$ – $88^{\circ}$ ,  $0^{\circ}$ E; (b)  $83^{\circ}$ – $88^{\circ}$ N,  $135^{\circ}$ E; (c)  $85.5^{\circ}$ – $86.5^{\circ}$ N,  $0^{\circ}$ E; and (d)  $86^{\circ}$ – $87^{\circ}$ N,  $135^{\circ}$ E. Troughs of this nature are pervasive on the polar caps of both hemispheres. Note that the scale shown here is exaggerated; the troughs are very shallow, with opening angles of up to  $165^{\circ}$ . Generally, the equator-facing wall is twice as steep as the opposing wall, Figure 7d being the most characteristic example.

probably warmer at depth because of geothermal heating. *Larsen and Dahl-Jensen* [2000] calculate that the basal temperature of the Martian north polar cap may be similar to that of terrestrial ice caps, that is, close to the freezing point, implying that flow speed will be governed primarily by deformation at the base and may be roughly the same as terrestrial ice caps of similar dimensions. Flow closure of the troughs, however, would be slower than terrestrial conditions since the near-surface ice in the cap will be closer to mean annual temperatures.

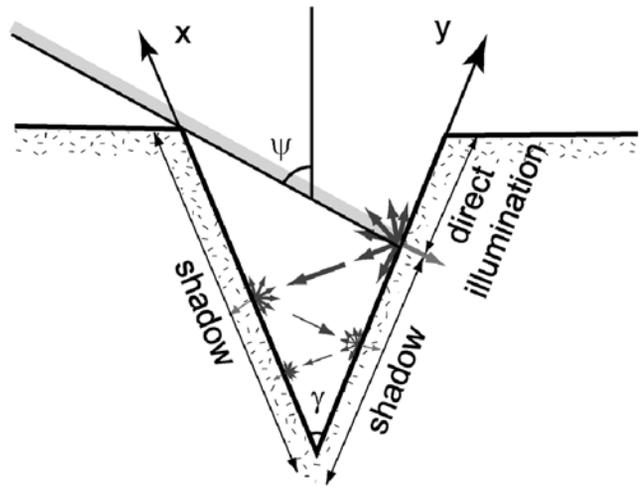
[25] Flow on the south polar cap is more difficult to assess since ice composition is still unknown. Pure  $\text{CO}_2$  ice is significantly softer than  $\text{H}_2\text{O}$  ice, which, combined with the overall cap geometry, suggests that the south polar cap cannot be pure  $\text{CO}_2$  [Nye, 2000].  $\text{H}_2\text{O}$ - $\text{CO}_2$  clathrate, however, is much stiffer than pure  $\text{H}_2\text{O}$  ice [Durham *et al.*, 2000]. At south polar thicknesses and temperatures a pure clathrate ice cap might flow so slowly as to be virtually immobile. Consequently, understanding the role of flow closure in the overall geometry of the south polar troughs will require significantly more study.

## 5. Proposed Evolution of the Martian North Polar Troughs

[26] The troughs cutting the north polar residual cap (Figure 8) vary greatly in depth. The shallowest troughs measured are  $\sim 100$  m deep; the deepest troughs are  $\sim 1$  km deep, cutting through the polar cap almost to the level of the surrounding plains. The troughs strike  $\sim 20^\circ$  north of west [Howard, 2000], which gives them their distinctive spiral look, and are steeper on their equator-facing walls than on their pole-facing walls [Ivanov and Muhleman, 2000b]. Equator-facing walls range from  $1^\circ$  to  $10^\circ$  slope; pole-facing walls range from  $<1^\circ$  to  $\sim 6^\circ$  slope. In general, wall slope seems to increase with wall height, so that the deepest grooves also have the steepest sides [Ivanov and Muhleman, 2000b]. Similarly, on the south polar residual cap, troughs range in depth from  $\sim 100$  m to  $>800$  m, with equator-facing walls ranging from  $1^\circ$  to  $9^\circ$  slope and pole-facing walls ranging from  $<1^\circ$  to nearly  $4^\circ$  slope [Ivanov and Muhleman, 2000b].

[27] Assuming the Martian ice cap troughs result from a series of events similar to those giving rise to the basins and canyons on the dry valley glaciers, we believe that the Martian troughs are primarily radiation-generated and radiation-maintained features. They are initiated high on the ice caps by some combination of small-scale surface roughness and dust. The small-scale roughness may be snow dunes [Howard, 2000] or something akin to the “cottage cheese” terrain photographed by MOC. This terrain is characterized by pitting of relatively regular width, varied length and connectivity, and depths of  $\sim 2$  to several meters and seems to cover most of the north polar residual cap [Thomas *et al.*, 2000; Malin *et al.*, 2001]. Either pitting or snow dunes will serve to trap dust, thereby lowering surface albedo and increasing energy receipt. Dune and pit geometry will also affect energy receipt as local slope and aspect increase or decrease incident radiation.

[28] As ice on the sides and floor of the dunes and pits sublimates, dust frozen into the ice will be freed. Accumulation of a thin layer of dust will increase ablation until the



**Figure 8.** Idealized V groove geometry used in the analysis. The  $x$  and  $y$  coordinates refer to nondimensional distances measured up the sides of the groove (both are 0 at the apex of the groove and 1 at the location where the groove wall intersects the flat surface into which the groove is cut). The angles  $\Psi$  and  $\gamma$  refer to the solar angle (described in the text) and the groove opening angle, respectively. The angle  $\Psi$  is related to the solar zenith angle. The enhancement of solar absorption due to multiple diffuse reflections of the incoming solar beam is shown schematically. Also shown are the portions of the groove that are shadowed from the incoming solar beam and the portions that are illuminated directly. The opening angle  $\gamma$  depicted here is much smaller than for the typical grooves on the northern polar ice sheet of Mars (see Table 2).

layer thickens and begins to insulate the ice. In a similar manner, ablation in the basins and canyons of the dry valley glaciers increases for thin debris cover and greatly decreases when the cover thickens sufficiently to insulate the ice. In the dry valleys, flowing water removes the debris from the ablation features. Current estimates of dust content of the Martian ice range from 50% to  $<1\%$  by volume [Clifford *et al.*, 2000]. This is sufficient to reduce the albedo and increase sublimation in the early stages of feature formation. However, sublimation of 1–10 m of ice will release ample dust to build up an insulating layer. Larger features will only form on the Martian residual caps in areas where wind transport can evacuate the excess dust. The wind will also aid in linking small basins and depressions into elongate features such as water links basins into canyons on Taylor Glacier. This suggests that mature Martian troughs form in geometries that optimize both radiation receipt and wind transport. This conclusion is supported by Howard [2000], who notes that the troughs strike at an angle  $20^\circ$  north of west, intermediate between the direction of greatest radiational loading and an orientation normal to the katabatic wind flow.

[29] The need for wind removal of the dust freed by ice sublimation may also explain the trough spacing on the ice caps. The troughs appear to be spaced fairly regularly at 20- to 70-km intervals. The ice between troughs forms a gentle bulge with a divide roughly halfway between, as can be seen in the MOLA data. We hypothesize that this is due to

**Table 2.** Parameters Used to Model the Martian Insolation Regime

Parameter	Value
Planetary distance, AU	1.524
Obliquity, deg	23.45
Length of day, hours	24.623
Planetary radius, km	3380
Solar constant, $W m^{-2}$	594
Opening angle, deg	165
Latitude, deg N	80

some channeling of the katabatic winds as they flow off the ice cap, possibly dictated in part by polar atmospheric waves. If the natural modes of the wind form some quasi-equally spaced banding, this, in turn, would favor dust deposition in or removal from roughness features at the center of those bands, leading to both faster development of roughness features in those regions and the linkage of those features downwind. This would further channel the wind, leading to a positive feedback loop and fairly regularly spaced mature troughs.

[30] Once initiated, the troughs will probably collect more debris on their floors than on their walls, leading to faster downward growth than lateral growth. As a result, the walls will steepen over time as the trough becomes deeper. This is supported by *Ivanov and Muhleman's* [2000a] measurements that indicate deeper troughs have steeper walls. Exposed dust banding in the equator-facing trough walls, seen in MOC imagery, indicates that the equator-facing walls experience net ablation. The pole-facing walls receive less radiation than the equator-facing walls as a result of their slope and aspect and will therefore be cooler. Sublimation on the pole-facing walls will be lower or nonexistent. MOC imagery suggests that the pole-facing walls may experience net accumulation since the debris banding in the ice seen on the equator-facing walls is not always visible on pole-facing walls.

## 6. Radiation Receipt Within the Martian North Polar Troughs: Why So Shallow?

[31] The enhancement of solar absorption derived from V groove geometry results from a geometrical “greenhouse” effect that stems from two physical principles. First, when the solar zenith angle is large, common at polar latitudes, an inclined surface oriented toward the equator gathers more sunlight. Therefore equator-facing slopes on terrestrial and Martian ice surfaces alike become steeper than opposing faces, which are shadowed, because they receive more energy and therefore have higher ablation rates. Second, outgoing reflected light from natural snow and ice surfaces can be treated as diffuse and varies with the cosine of the incidence angle of the incoming beam. A portion of the diffuse reflected energy is directed toward the opposing wall where it is reflected again. The multiplicity of reflections associated with the diffuse radiation leads to greater absorption because some fraction of the incident energy is absorbed with each reflection (Figure 8).

[32] The radiative regime of an idealized Martian trough with perfect “V” geometry is computed following *Pfeffer and Bretherton's* [1987] work on crevasses on terrestrial glaciers acting as solar radiation traps. We chose parameters

to represent typical valleys on the Martian northern ice sheet (Table 2), which are derived from a visual analysis of the MOLA transects and valley profiles. We ignored secondary effects such as angular heterogeneity in the absorptivity of an icy surface and the effects of “far-field” topographic patterns, such as the general slope of the ice sheet and the effects of distant ice structures that may cast shadows or radiate energy on points within the groove. Small-scale roughness on the sides of the groove, such as the “stair step” topography found in the dry valleys, is also not considered, though there is some evidence in the MOC imagery that stair step topography may occur within the Martian grooves.

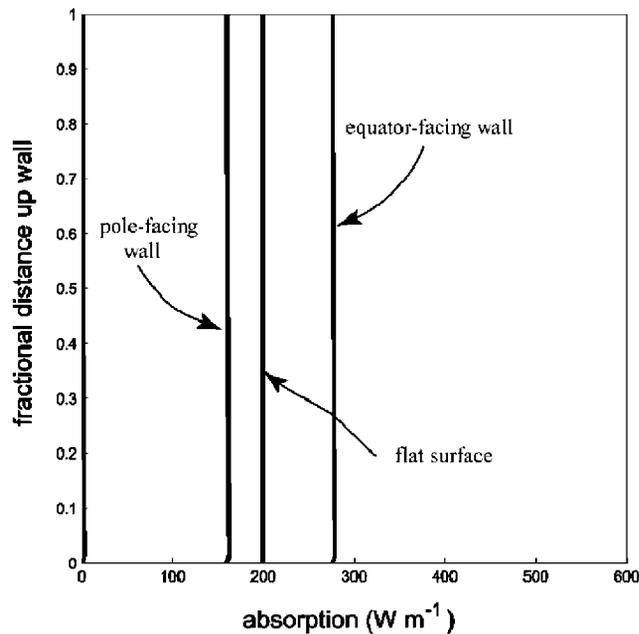
[33] As illustrated in Figure 8, variables  $x$  and  $y$  are nondimensional coordinates that describe distance along each of the two walls from an origin at the bottom of the V. The  $y$  coordinate is associated with the wall that faces the Sun, the equator-facing wall of the groove. The coordinate axes are separated by the angle  $\gamma$ , the V groove opening angle. Thus the coordinate axes are not Cartesian. The variable  $\Psi$  is the angle between the direct solar beam and the line that bisects  $\gamma$ . The angle  $\Psi$  is related to the solar zenith angle, which depends on the time of year and the time of day, as well as other long-term variables such as the planet's obliquity. The variable  $\Psi$  also depends on the inclination of the V groove's line of symmetry from the vertical. For typical grooves on the ice sheets of Mars the equator-facing wall is steeper than the opposing wall by about a factor of 2 and ranges in slope from  $4^\circ$  to  $10^\circ$ . Thus the line of symmetry is typically inclined toward the equator by  $1^\circ$ – $4^\circ$ .

## 7. Results and Discussion

[34] Figure 9 shows the results for a typical Martian groove, where  $\gamma \sim 165^\circ$  (the equator-facing wall is at a  $10^\circ$  slope, and the pole-facing wall is at a  $5^\circ$  slope) and where englacial and surficial debris causes the reflectivity to be much reduced (we use  $\rho = 0.3$  on the equator-facing wall and  $\rho = 0.4$  on the pole-facing wall). The calculation displayed in Figure 9 represents the circumstances encountered at  $80^\circ N$  at noon on the summer solstice (for a groove that extends east to west).

[35] The most important result, shown in Figure 9, is that energy absorption on the Sun-facing wall is significantly enhanced over the energy absorption on a flat, horizontal surface. The Sun-facing wall, because it is inclined toward the incoming solar beam, absorbs  $\sim 140\%$  of the energy that a flat, horizontal surface would absorb. The enhancement of incoming energy is analogous to that observed for the basins and canyons on the Antarctic dry valley glaciers. Reflected radiation accounts for 2% of the total absorbed radiation. The Sun-opposing wall absorbs  $\sim 80\%$  of that on a flat, horizontal surface. On the Sun-facing wall, reflected radiation accounts for  $<1\%$  of the total.

[36] The low values for reflected radiation imply that the large opening angles of the Martian grooves do not allow the groove walls to “see” each other. This implies that reradiated longwave radiation absorption will also be low. These results show that a relatively simple parameterization of incoming radiation can be used for the grooves, such as the straight cosine dependence of *Ivanov and Muhleman*



**Figure 9.** Absorption by the V groove geometry using  $\gamma = 165^\circ$  at local noon on the summer solstice at  $80^\circ\text{N}$ . See also Figures 2 and 3 of Pfeffer and Bretherton [1987] for further details and results applicable to terrestrial glaciers. Curves indicate the absorbed energy for the Sun-facing and Sun-opposing walls and for a flat surface.

[2000a, equation (6)]. Omitting multiple reflections only incurs an error of  $\sim 2\%$  in the radiation calculations.

[37] The radiation regime is computed for  $\gamma = 60^\circ$ , the approximate opening angle of the Taylor Glacier canyons (the equator-facing wall is vertical, and the pole-facing wall is at  $30^\circ$ ), to show the sensitivity to the groove opening angle (Figure 10). At this opening angle the pole-facing wall and the lower portion of the equator-facing wall are both shadowed. The lower portion of the equator-facing wall absorbs up to  $20 \text{ W m}^{-2}$  of shortwave radiation because of reflection off the pole-facing wall. The pole-facing wall receives up to  $100 \text{ W m}^{-2}$  from reflection off the equator-facing wall. Comparison with Figure 9 shows that as  $\gamma$  is reduced, the overall efficiency of the groove in capturing solar radiation becomes greater relative to that of a flat, horizontal surface. The incident radiation on the equator-facing wall is strongly enhanced because the wall faces more directly into the Sun, and the reflected radiation component of the total absorbed radiation is large because of the narrow opening angle of the groove.

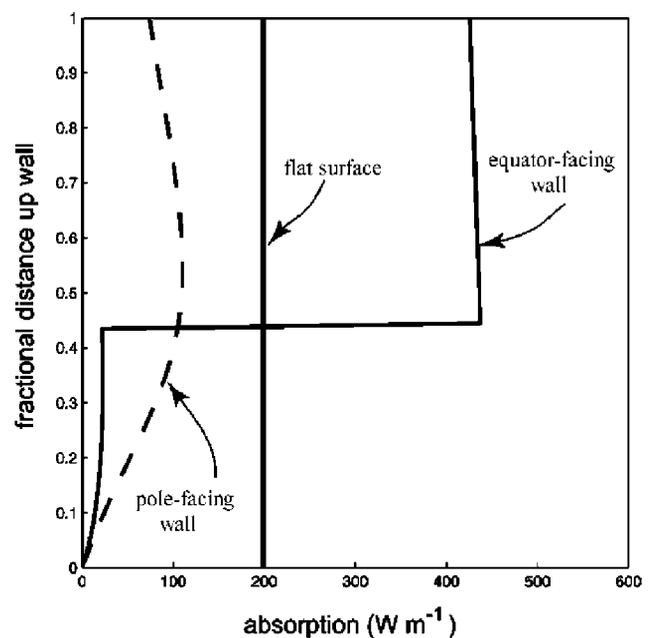
[38] The variation of absorption with opening angle prompts us to speculate on the possible time evolution of Martian grooves. In model runs for grooves with  $\gamma < \sim 60^\circ$  the energy is strongly concentrated near the top of the groove. If allowed to generate ablation through sublimation, this energy concentration will cause the groove to widen near its top and thus cause  $\gamma$  to increase. In model runs with  $\gamma > \sim 90^\circ$  the radiation absorbed by the two walls becomes more or less uniform over the length of the walls, as is the case displayed in Figure 9. Further sublimation is then expected to enlarge the groove, maintaining its depth to width ratio, and tilt its axis of symmetry toward the Sun.

Thus an opening angle of  $\sim 90^\circ$  or more appears to be “optimum” in the sense that once it is achieved, the groove will both deepen and widen while maintaining a constant shape (i.e., a constant  $\gamma$ ).

[39] However, if  $\gamma = \sim 90^\circ$  is “optimum,” why do the troughs on the Martian residual caps have opening angles of  $165^\circ$ ? On Taylor Glacier the canyons can be characterized as grooves with an opening angle of  $60^\circ$ – $90^\circ$ . In particular, the pole-facing walls slope quite consistently  $30^\circ$ – $35^\circ$ , an angle that serves to minimize the incoming direct solar beam. This is clearly not the case on the Martian ice caps where the pole-facing walls slope at angles of  $1^\circ$ – $6^\circ$ . Similarly, the slopes seen on the equator-facing walls do not maximize incoming radiation, unlike in the dry valleys of Antarctica where equator-facing slopes are near vertical. This suggests that on Mars some mechanism, such as closure by flow, is at work to keep the slopes relatively shallow. If so, quantifying radiation receipt and sublimation within the troughs may allow us to estimate ice flow rates.

## 8. Some Speculations on the Origin of the “Swiss Cheese” Terrain

[40] Thomas *et al.* [2000], Byrne and Ingersoll [2000], and Malin *et al.* [2001] attribute the basins in the south polar “Swiss cheese” terrain to sublimation. This would explain their circular shape and steep walls and the appearance in MOC imagery that they grow outward from an initial point, pass through a stage where they overlap in places and leave oddly shaped mesas in others, and ultimately leave behind a new surface, flattened at the level of the basin floors. However, if sublimation and high latitude were the only



**Figure 10.** Absorption by the V groove geometry using  $\gamma = 60^\circ$  at local noon on the summer solstice at  $80^\circ\text{N}$ . The curves indicate the absorbed energy for the equator-facing and pole-facing walls and for a flat surface. The equator-facing wall is shadowed from the bottom of the groove to about  $y = 0.44$ .

factors required for basin growth, we should also see Swiss cheese terrain on the north polar cap. Since we do not, the presence of these basins only on the south polar residual cap must relate to the physical differences between the caps. Probably this difference relates to the presence of dry ice in the south polar cap and not in the north.

[41] We propose that the flat floors of the basins and the fact that the basins appear to grow together only through radial growth is indicative of both an environment where wind transport of debris is low and/or where basin formation is rapid. As a result, we believe the basins indicate that the ice in the region where they are found is CO<sub>2</sub> ice, a medium that can rapidly ablate under south polar residual cap summer conditions.

[42] The Swiss cheese terrain is visually similar to the topography of Canada Glacier where the basins are fairly randomly distributed (Figure 4). Unlike the canyons seen on Taylor Glacier, which are initiated by albedo differences resulting from the presence of medial moraines and which therefore form features in linear arrays that are subsequently linked by water transport of debris, the Canada Glacier basins have no underlying linking mechanism. Similarly, we believe that unlike the Martian troughs, which we have proposed require wind removal of debris both to prevent lag deposits from insulating the ice, preventing ablation, and to link smaller features into elongated grooves, the Martian Swiss cheese basins require fast growth or limited wind influence to explain why they are not linked into trough features.

[43] If the Swiss cheese basins form in CO<sub>2</sub> ice, the basins walls could potentially ablate back on the order of a meter over the course of the summer (1–2 m of CO<sub>2</sub> frost sublimates off the south polar cap during the summer). Erosional widening of pits has been observed at the predicted rates [Malin *et al.*, 2001], adding further support to inferences that these terrains are made of dry ice. Sublimation rates on the near-vertical walls of the Swiss cheese basins and on deposits at the foot of the scarps should be higher than sublimation from horizontal surfaces because near-vertical walls will capture more radiation, as discussed for the north polar cap in section 6. The debris contained in this ice would be deposited at the base of the wall, which could further enhance sublimation rates, thus creating moats, as observed near some scarps. However, this does not explain why fields of basins appear to maintain a relatively consistent depth. In the dry valleys, neighboring basins can be of very different depths. This suggests that the depths of the Martian basins are controlled by an external mechanism such as a change in ice composition.

## 9. Conclusions

[44] The dry valley glaciers in Antarctica appear to serve as a useful analogue for the Martian polar caps. Study of the basins and canyons seen on many dry valley glaciers provides us with insight into the sequence of events necessary for the initiation and growth of large-scale features. Studying the dry valley glaciers also highlights the general processes required to maintain such features once they have formed. In particular, such comparison highlights the requirement of some mechanism to remove debris from the ice surface. Without such a mechanism the debris can

rapidly accumulate and insulate the ice from further sublimation. This implies that on Mars, there must be eolian removal of dust freed by sublimation from trough walls, which, in turn, indicates that the troughs are the result of some radiation and wind interaction.

[45] Using the *Pfeffer and Bretherton* [1987] V groove radiation model, adapted for Mars, our results indicate that although the grooves on Martian ice sheets have a very large opening angle, there is still a significant increase in radiation absorption of the equator-facing wall. What is surprising, however, is that the large-scale grooves on the ice sheets have opening angles that are too wide to optimize solar absorption. We speculate that the reason for this inefficiency is that groove closure by ice flow overpowers solar energy-driven sublimation. This suggests that further work aimed at quantifying the energy absorption and sublimation within the grooves may allow us to estimate ice flow rates involved in the flow closure of the grooves, a topic we hope to address in future work.

[46] Comparison of the Canada Glacier basins and the Martian south polar “Swiss cheese” terrain suggests that sublimation could be responsible for excavating the circular depressions seen. We believe that the random distribution of the Swiss cheese basins indicates that wind removal of debris does not play a role in their development. We further suggest that this indicates that the basins form rapidly, perhaps indicating that the ice in which they are found is CO<sub>2</sub> ice. However, the consistent depth of basins over large areas indicates some mechanism of depth control such as a layer of “harder,” less volatile H<sub>2</sub>O or CO<sub>2</sub> clathrate ice.

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